Quantum Coherence between Q-bits and Thermal Environments: Non-Markovian Approach.

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Systems in Thermal Environments



Standard Thermodynamics:

Weak Coupling between System (S) and Environment (B)

- Usually, it is strong for small systems.
- Even if it is weak, Entanglement can be formed.

Open Quantum Systems: Exact Scenario

Hamiltonian:
$$H_{\rm SB} = H_{\rm S} + H_{\rm B} + V_{\rm SB}$$

Unitary Evolution $i \frac{\partial \rho_{\rm SB}}{\partial t} = [H_{\rm SB}, \rho_{\rm SB}]$

State of the System $ho_{
m S}(t) = {
m tr}_{
m B}[
ho_{
m SB}(t)]$

S: System, B: Thermal Bath, S+B: Isolated System

- Practically Impossible
- Need an equation of motion for $ho_{
 m S}(t)$

Open Quantum Systems: Born-Markovian Approximation

Weak Coupling between S and B: Born Approximation Short correlation time for B: Markovian Approximation

Quantum Master Equation

In energy eigenbasis,

- Off-diagonal element vanishes very quickly (Decoherence)
- The transition between eigenstates is incoherent.
- No entanglement between S and B

Composite Systems



Four time scales 1)system relaxation time 2)memory in environment 3)S-B coupling strength 4)Internal coupling strength



$$B_1 \cdots \boxed{\blacksquare} \cdots$$

Specific Model

Coupled Q-bits

$$\hat{H}_{S} = \frac{\omega_{0}}{2}\hat{\sigma}_{S_{1}}^{z} + \frac{\omega_{0}}{2}\hat{\sigma}_{S_{2}}^{z} + \lambda_{S}\left(\hat{\sigma}_{S_{1}}^{+}\hat{\sigma}_{S_{2}}^{-} + \hat{\sigma}_{S_{1}}^{-}\hat{\sigma}_{S_{2}}^{+}\right)$$

$$E_{1} = -\lambda_{S}, \quad |u_{1}\rangle = |--\rangle$$

$$E_{2} = -\omega_{0}, \quad |u_{2}\rangle = \frac{1}{\sqrt{2}}\left(|+-\rangle - |-+\rangle\right)$$

$$E_{3} = +\omega_{0}, \quad |u_{3}\rangle = \frac{1}{\sqrt{2}}\left(|+-\rangle + |-+\rangle\right)$$

$$E_{4} = +\lambda_{S}, \quad |u_{4}\rangle = |++\rangle$$

Boson Baths

$$\hat{H}_{B_i} = \sum_k \omega_{B_i}(k) \,\hat{a}_{B_i}^{\dagger}(k) \hat{a}_{B_i}(k), \qquad i = 1, 2$$

System-Bath Coupling

$$\hat{V}_{S_i B_i} = \hat{X}_{S_i} \otimes \hat{Y}_{B_i}$$
$$\hat{X}_{S_i} = \sigma_{S_i}^+ + \sigma_{S_i}^-$$
$$\hat{Y}_{B_i} = \sum_k \epsilon_{B_i}(k) \left[\hat{a}_{B_i}^\dagger(k) + \hat{a}_{B_i} \right]$$

$$B_1 \cdots \textcircled{3} B_2$$



Drude-Lorenzian model

$$g_{B_i}(\omega) = \sum_k |\epsilon_{B_i}(k)| \delta \left(\omega - \omega_{B_i}(k)\right)$$
$$= \frac{2\lambda_{B_i}\gamma_{B_i}\omega}{\omega^2 + \gamma_{B_i}^2}$$

Unitary Evolution of the total system

$$\rho_{SB}(t) = \left\{ \overleftarrow{T} \prod_{i} e^{-i \int_{t_0}^t \hat{V}_{S_i B_i}(s) \mathrm{d}s} \right\} \rho_{SB}(t_0) \left\{ \overrightarrow{T} \prod_{i} e^{i \int_{t_0}^t \hat{V}_{S_i B_i}(s) \mathrm{d}s} \right\}$$

Initial state: $\rho_{SB}(t_0) = \rho_S(t_0) \otimes \rho_{B1}(t_0) \otimes \rho_{B_2}(t_0)$ Wick's theorem works: $\rho_{B_i}(t_0)$ is a quasi-free state. $\rho_{B_i}(t_0) = \frac{1}{Z_i} e^{-\beta_i \hat{H}_{B_i}}$

$$\rho_S(t) = \operatorname{tr}_{B_i}\left[\rho_{SB}(t)\right] = \overleftarrow{\mathfrak{T}} \prod_i e^{-\int_{t_0}^t \int_{t_0}^{t_1} \mathrm{d}t_1 \mathrm{d}t_2 \mathcal{K}_i(t_1, t_2)} \rho_S(t_0)$$

$$\begin{aligned} \mathcal{K}_{i}(t_{1}, t_{2}) &= S_{i}^{-}(t_{1}) \operatorname{Im} C_{i}(t_{1} - t_{2}) S_{i}^{-}(t_{2}) + i S_{i}^{-}(t_{1}) \operatorname{Re} C_{i}(t_{1} - t_{2}), S_{i}^{+}(t_{2}) \\ S_{i}^{\pm}(t) &= \left[\hat{X}_{S_{i}}(t), \cdot \right]_{\pm} \\ C_{i}(t) &= \left\langle \hat{Y}_{B_{i}}(t) \hat{Y}_{B_{i}}(t_{0}) \right\rangle_{0} \approx \lambda_{B_{i}} \left[c_{i} e^{-\gamma_{B_{i}} t} + 2\Delta_{i} \delta(t) \right] \\ c_{i} &= 2/\beta_{i} - \gamma_{B_{i}}(\Delta_{i} + i) \qquad \Delta_{i} = B_{i} \beta_{i}/6 \end{aligned}$$

Auxiliary operators

$$\sigma_{n_1,n_2}(t) = \overleftarrow{\mathfrak{T}} \prod_i \left\{ \left[-i \int_{t_0}^t \mathrm{d}s \, e^{-\gamma_{B_i}(t-s)} \mathcal{G}_i(s) \right]^{n_i} \\ \times e^{-\lambda_{B_i} \int_{t_0}^t \int_{t_0}^{t_1} \mathrm{d}t_1 \mathrm{d}t_2 \mathcal{S}_i^-(t_1) e^{-\gamma_{B_i}(t_1-t_2)} \mathcal{G}_i(t_2)} \\ \times e^{-\lambda_{B_i} \Delta_i \int_{t_0}^t \mathrm{d}t_1 \mathcal{S}^-(t_1) \mathcal{S}^-(t_1)} \right\} \rho_S(t_0)$$

$$\mathcal{G}_i(t) = \left(2/\beta_i - \gamma_{B_i} \Delta_i\right) \mathcal{S}^-(t) - i\gamma_{B_i} \mathcal{S}^+(t).$$



Equation of Motion

$$\frac{\mathrm{d}}{\mathrm{d}t}\sigma_{n_{1},n_{2}}(t) = -(\gamma_{B_{1}}n_{1} + \gamma_{B_{2}}n_{2})\sigma_{n_{1},n_{2}}(t)
- \left[\lambda_{B_{1}}\Delta_{1}S_{1}^{-}(t)S_{1}^{-}(t) + \lambda_{B_{2}}\Delta_{2}S_{2}^{-}(t)S_{2}^{-}t)\right]\sigma_{n_{1},n_{2}}(t)
- i\lambda_{B_{1}}S_{1}^{-}\sigma_{n_{1}+1,n_{2}}(t) - i\lambda_{B_{2}}S_{2}^{-}\sigma_{n_{1},n_{2}+1}(t)
- in_{1}\lambda_{B_{1}}\mathcal{G}_{1}(t)\sigma_{n_{1}-1,n_{2}}(t) - in_{2}\lambda_{B_{2}}\mathcal{G}_{2}(t)\sigma_{n_{1},n_{2}-1}(t)$$

HEOM is in principle exact.

- Correct description of coherence within the system
- Coherence with environment including entanglement.
- Strong coupling.

Heat Current

$$J_{i} = -\frac{\mathrm{d}}{\mathrm{d}t} \langle H_{B_{i}} \rangle = -i \operatorname{tr}_{S} \left\{ [\hat{X}_{S_{i}}, \eta_{i}] \hat{H}_{S} \right\}$$
$$\eta_{1} = \lambda_{B_{1}} \left\{ \sigma_{1,0} - i \Delta_{1} \mathcal{S}_{1}^{-} \sigma_{0,0} \right\}$$
$$\eta_{2} = \lambda_{B_{2}} \left\{ \sigma_{0,1} - i \Delta_{2} \mathcal{S}_{2}^{-} \sigma_{0,0} \right\}$$

Numerical Calculation Procedure (Exact enumeration in principle)

- Initial condition: $\sigma_{0,0}(t_0) = \rho_S(t_0), \quad \sigma_{i,j}(t_0) = 0$
- Hierarchic cut-off: $\sigma_{i,j}(t) = 0$ for $i > i_{\max}, j > j_{\max}$
- Solve coupled ODEs (HEOM) numerically
- Non-Markovian solution: $ho_S(t) = \sigma_{0,0}(t)$



 $\omega_0 = 1, \quad \lambda_S = 0.5$ $\lambda_{B_1} = \lambda_{B_2} = 0.1$ $\gamma_{B_1} = 0.2, \quad T_1 = 2$ $\gamma_{B_2} = 0.1, \quad T_2 = 1$

Hierarchy depth = 20

Comparison between QME and HEOM

Heat Current



Comparison between QME and HEOM

Entanglement between Q-bits





Trace distance =
$$\|\rho - \sigma\| = \frac{1}{2} \operatorname{tr} \left[\sqrt{(\rho - \sigma)^{\dagger} (\rho - \sigma)} \right]$$

Autonomous Quantum Heat Engine



